

## Agenda of July 21, 2026 zoom meeting

1. Matt      5' Status of PRD / arXiv paper submission
2. Carolyn Haviland DPF talk at FNAL will be posted at INDICO for comments. Her actual presentation is on Thursday.
3. Yuri      10' Possible physics outcomes of  $n$ TTM-2 experiment
4. Yuri      15' Results from mu-shield prototype at UT
5. Mubi      5' Updates from UKY
6. Linus      5' Updates from LU
7. John      15' Discussion of Design Review questions for tomorrow

News: Nathalie Ziehl (PSI) thesis with recent UCN  $n \rightarrow n'$  search limits can be found at <https://doi.org/10.3929/ethz-c-000800962>.

Zsigmond Geza (PSI) reconfirmed that they will send us soon digitally their  $\tau$  vs  $\Delta m$  limits

# Possible physics outcomes of $n$ TTM-2 experiment

Yuri Kamyshev

In the framework of evolution of the mixed  $(n \ n')^T$  system

$$i\hbar \begin{pmatrix} \dot{n} \\ \dot{n}' \end{pmatrix} = \mathcal{H} \begin{pmatrix} n \\ n' \end{pmatrix}$$

we can consider different “magnetically compensated” Hamiltonians

$$\mathcal{H} = \begin{pmatrix} U - \mu B & \delta \\ \delta & 0 \end{pmatrix}$$

that can be explored with regeneration  $n$ TMM experiment.

In other words, we can analyze  $n$ TMM data in different theoretical models.  
Or in the absence of resonance peak set limits on different parameters.

Assume two 30-min scans with  $\sim 80$  points each with step 0.04 G

# 1. $n$ TMM original scheme

$$\mathcal{H} = \begin{pmatrix} V_F - \mu B & \epsilon + \kappa \mu B \\ \epsilon + \kappa \mu B & 0 \end{pmatrix} \quad \epsilon \ll \kappa \mu B$$

Procedure: scan only around  $V_F - \mu B = 0$  is needed

Result : value of  $n$ TMM  $\kappa$

## 2. $n$ TMM $\kappa$ and $\pm\Delta m$

$$\mathcal{H} = \begin{pmatrix} V_F \pm \Delta m - \mu B & \epsilon + \kappa\mu B \\ \epsilon + \kappa\mu B & 0 \end{pmatrix} \quad \epsilon \ll \kappa\mu B$$

Procedure: full scan in the range  $= 2\Delta m$

Result : 2-dimensional limits of

$\kappa$  value and  $\Delta m \in \pm 0.012 \text{ neV } (\pm 2 \text{ G})$

3. no  $n$ TMM, no  $\Delta m$  - only  $\epsilon$

$$\mathcal{H} = \begin{pmatrix} V_F - \mu B & \epsilon \\ \epsilon & 0 \end{pmatrix}$$

Procedure: scan only around  $V_F - \mu B = 0$

Result : limit on value of  $\epsilon$  ( $\Delta m = 0, B' = 0$ )

Likely not very competitive (except by cold beam)

4. no  $n$ TMM, but  $\epsilon$  and  $\pm \Delta m$

$$\mathcal{H} = \begin{pmatrix} V_F \pm \Delta m - \mu B & \epsilon \\ \epsilon & 0 \end{pmatrix}$$

Procedure: full scan range =  $2\Delta m$

Result : 2-dimensional limits of

$\epsilon$  vs  $\Delta m \in \pm 0.012 \text{ neV } (\pm 2 \text{ G})$

Likely not very competitive (except by cold beam)

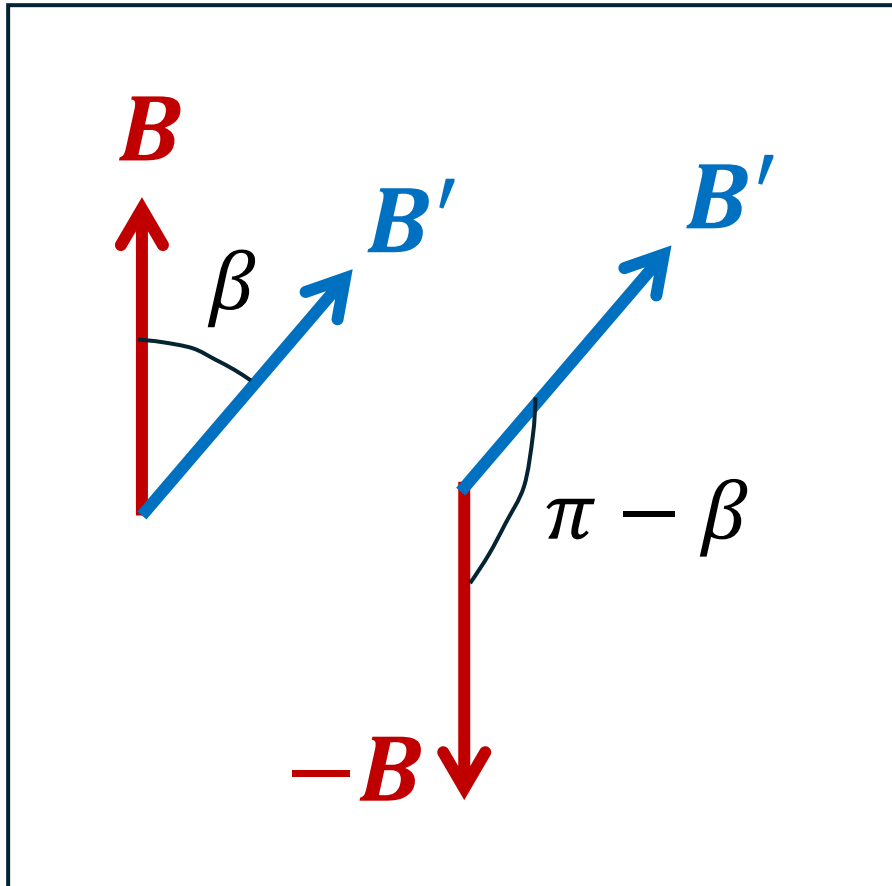
## 5. Mirror field $B'$ and $\epsilon$ ; no $n$ TMM

$$\mathcal{H} = \begin{pmatrix} V_F + \mu \vec{\sigma} \cdot \vec{B} & \epsilon \\ \epsilon & \mu \vec{\sigma} \cdot \vec{B}' \end{pmatrix}$$

Procedure: Full range scan  $\sim \pm 2G$  (to compensate  $B'$ )

Result : 2-dim limits  $\epsilon$  (or  $\tau$ ) vs  $B'$

- Width of expected resonance is a QM effect depending on the average TOF provided by  $n$ -spectrum (similar for all)
- We might have an advantage of cold beam vs UCN in  $B'$  search. In our cold beam scenario  $B'$  signal doesn't depend on  $\beta$ . (assuming that depolarization in regeneration is not essential)



$$P_B(t) = \mathcal{P}_B(t) + \mathcal{D}_B(t) = \mathcal{P}_B(t) + D_B(t) \cos \beta,$$

where  $\beta$  is the angle between the vectors  $B$  and  $B'$  and

$$\mathcal{P}_B(t) = \frac{\sin^2[(\omega - \omega')t]}{2\tau^2(\omega - \omega')^2} + \frac{\sin^2[(\omega + \omega')t]}{2\tau^2(\omega + \omega')^2},$$

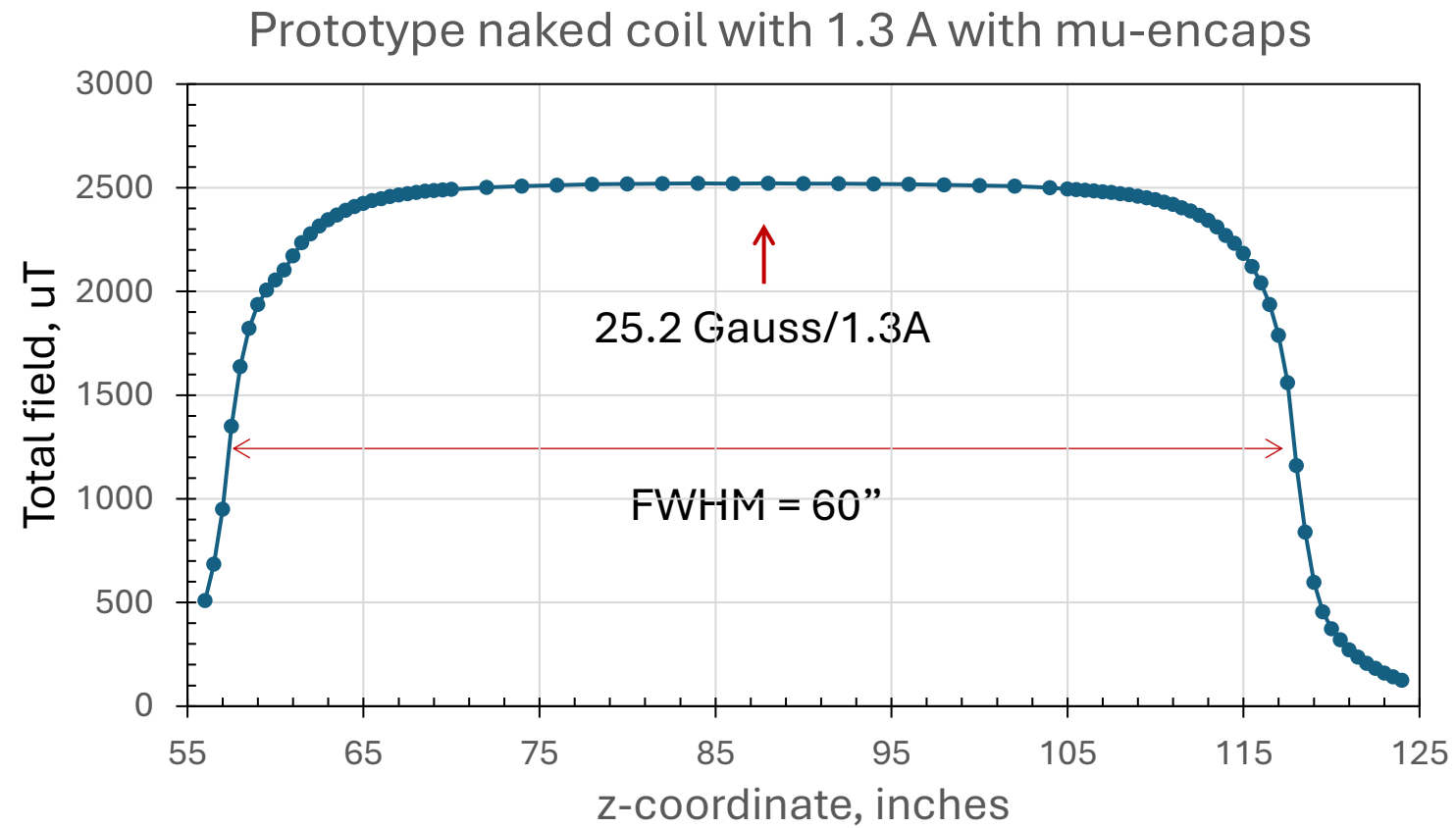
$$D_B(t) = \frac{\sin^2[(\omega - \omega')t]}{2\tau^2(\omega - \omega')^2} - \frac{\sin^2[(\omega + \omega')t]}{2\tau^2(\omega + \omega')^2},$$

$$\text{with } \omega = \frac{1}{2}|\mu B| \text{ and } \omega' = \frac{1}{2}|\mu B'|. \quad \tau = \hbar/\epsilon$$

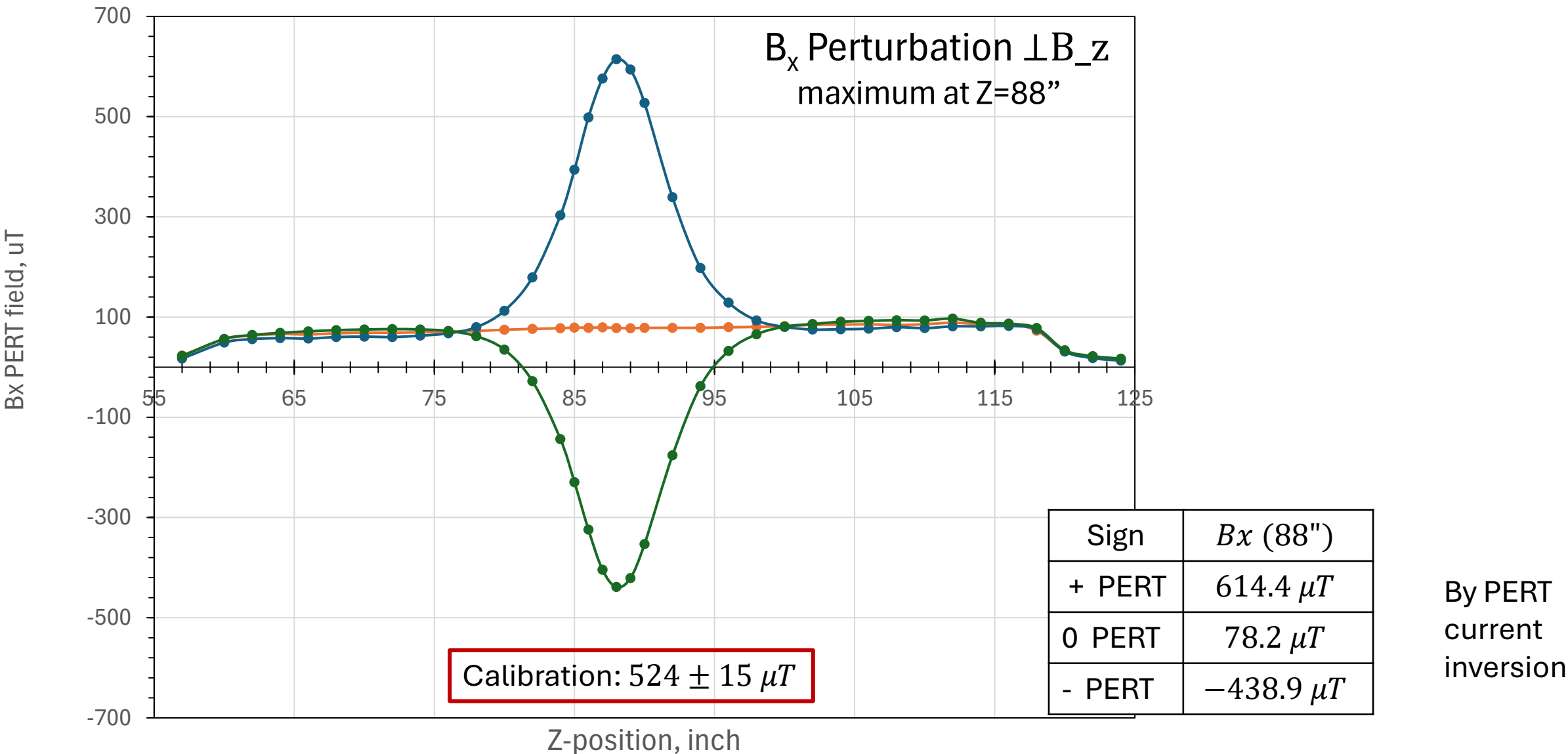
# Results from $\mu$ – shield prototype at UT

Y. Kamyshev & S. Vavra

# Naked Coil, no mu-metal, no PERT

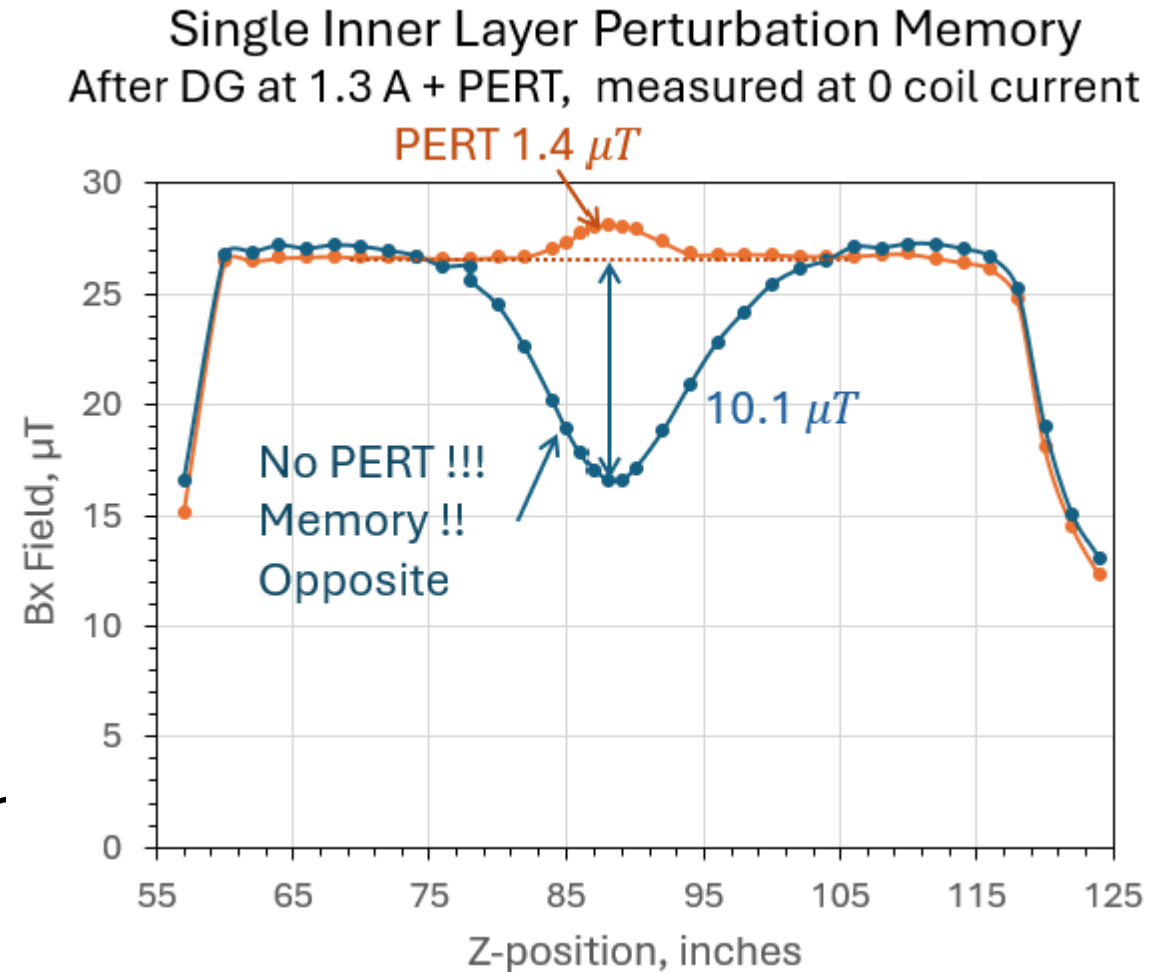


# Naked Coil, with PERT, Coil current 1.3 A



- Degaussing of mu-metal should be performed with the magnet and perturbation state close to operational !

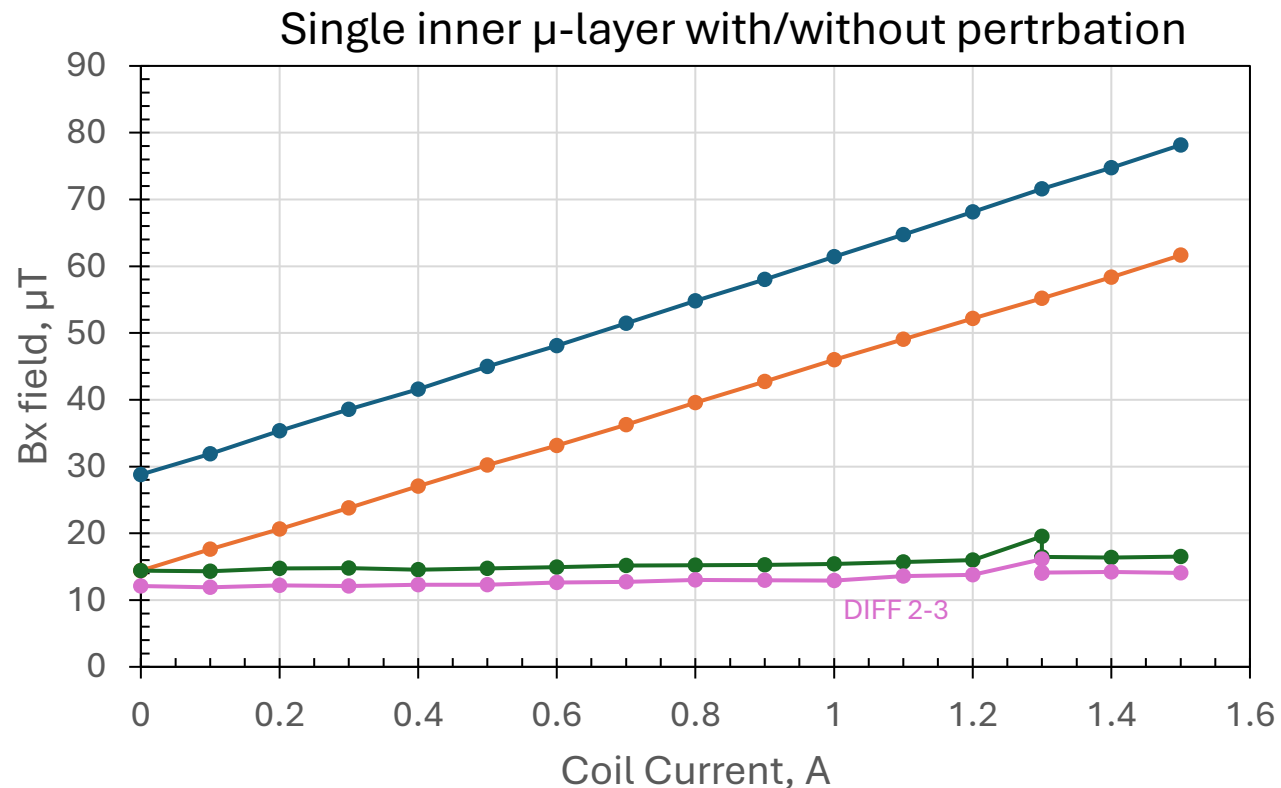
- Our standard mode of degaussing was with + PERT and coil 1.3 A
- By degaussing the shielding remembers the perturbation with opposite sign →
- For Shielding Factor measurement it was important to set magnetometer position on the peak of PERT



# 1 $\mu$ - inner layers

- Major question was how shielding factor will change with the saturation of mu-metal serving as flux return yoke of the solenoid.

Standard degaussing,  $B_x$  field with/without perturbation at z peak position 88"

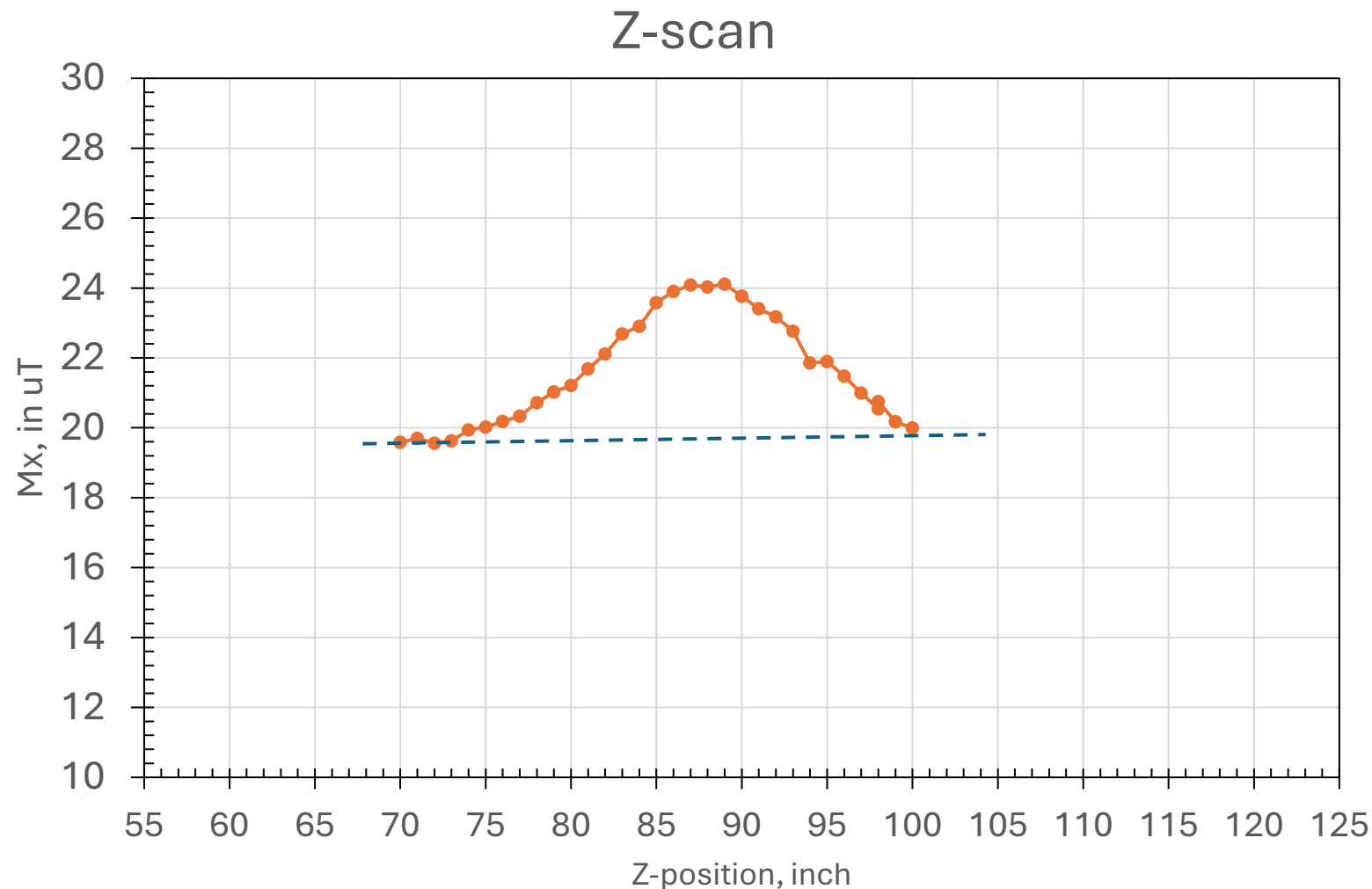


Single inner mu-layer  
Seen PERT at 1.3 A  
 $\overline{\text{PERT}} = 16.1 \pm 1.8 \mu T$

DIFF dependence on current  
 $\sim 14.18 + 1.47 \cdot i(A) \mu T$

## 2 $\mu$ - inner layers

Double shielding is good. Difficult to see the perturbation peak after 2- $\mu$  shield  
Special degaussing with 0A and NO PERT. Then, Z-scan with PERT, current 0 A



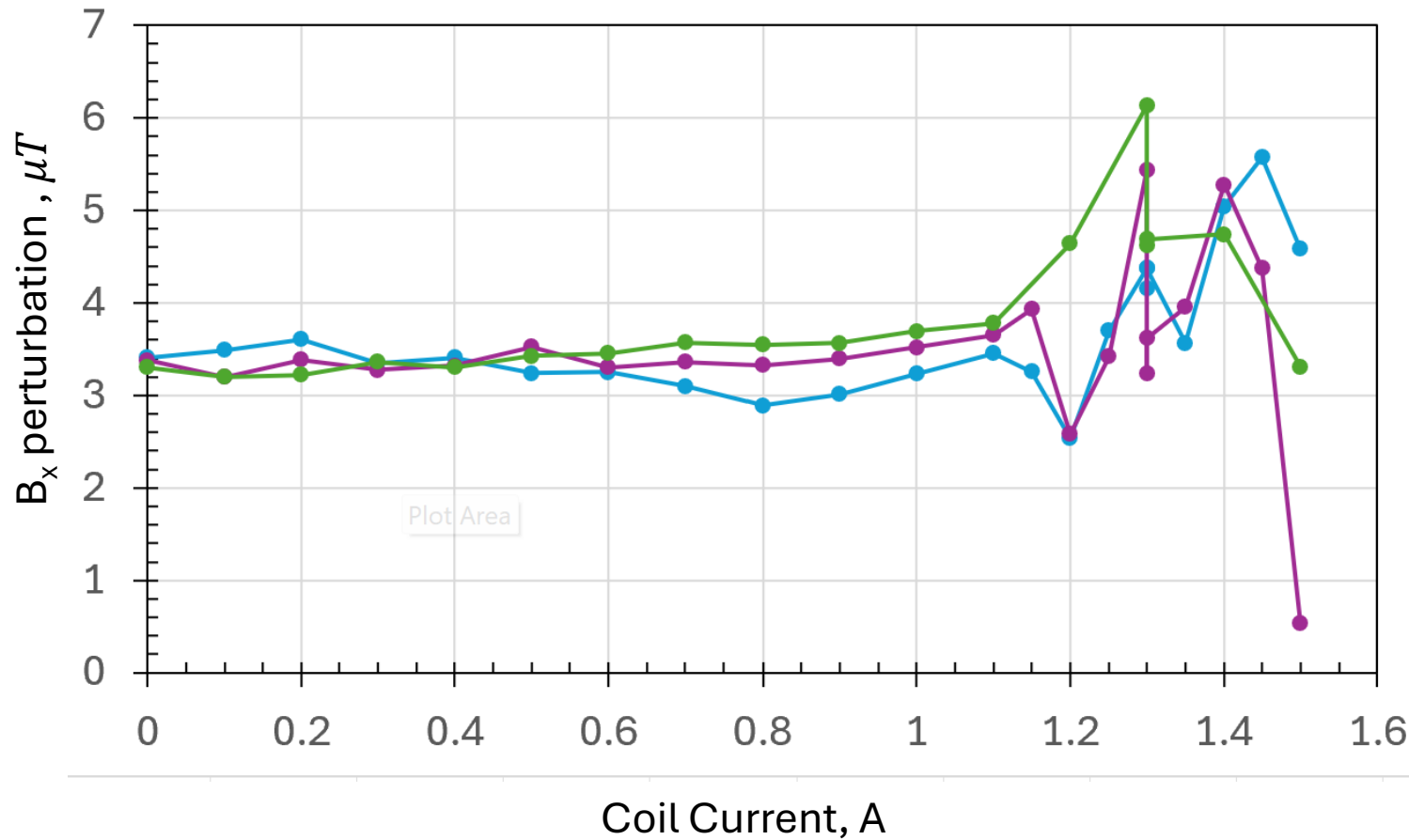
Peak position 88"  
Perturbation  $\sim 4.5 \mu T$

## 2 $\mu$ - inner layers

After Standard Degaussing, Z=88". DIFF=PERT-NOPERT, with Air

At smallest PERT signals (double layer) with large coil current temp. regime is not stable

Perturbation Difference vs Coil Current



Double inner mu-layer

Seen PERT at 1.3 A

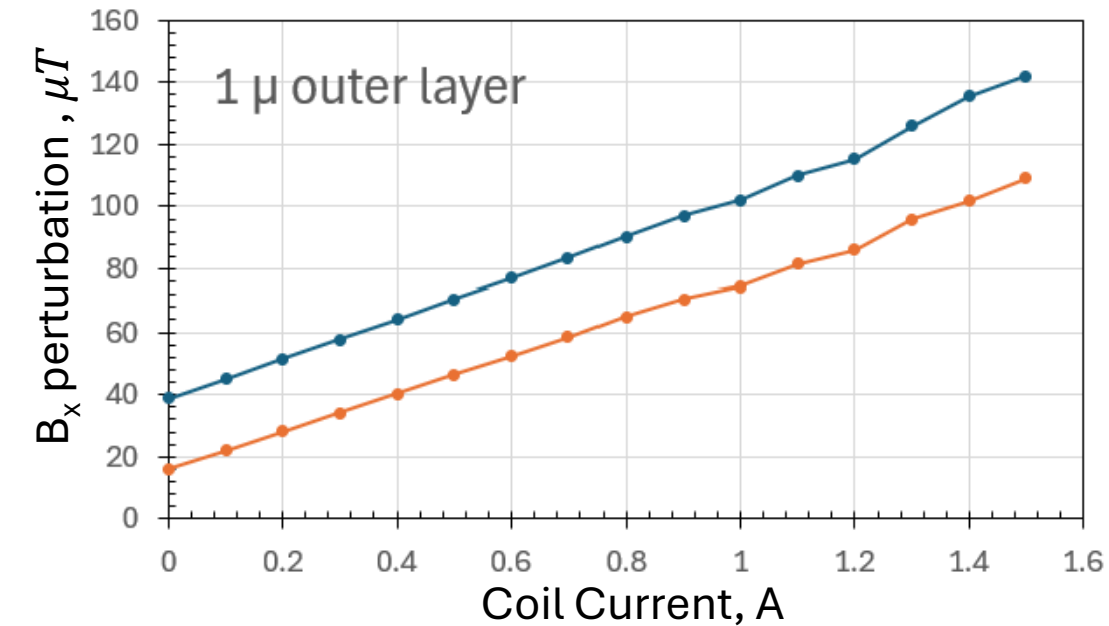
$$\overline{\text{PERT}} = 4.51 \pm 0.61 \mu T$$

DIFF dependence on current

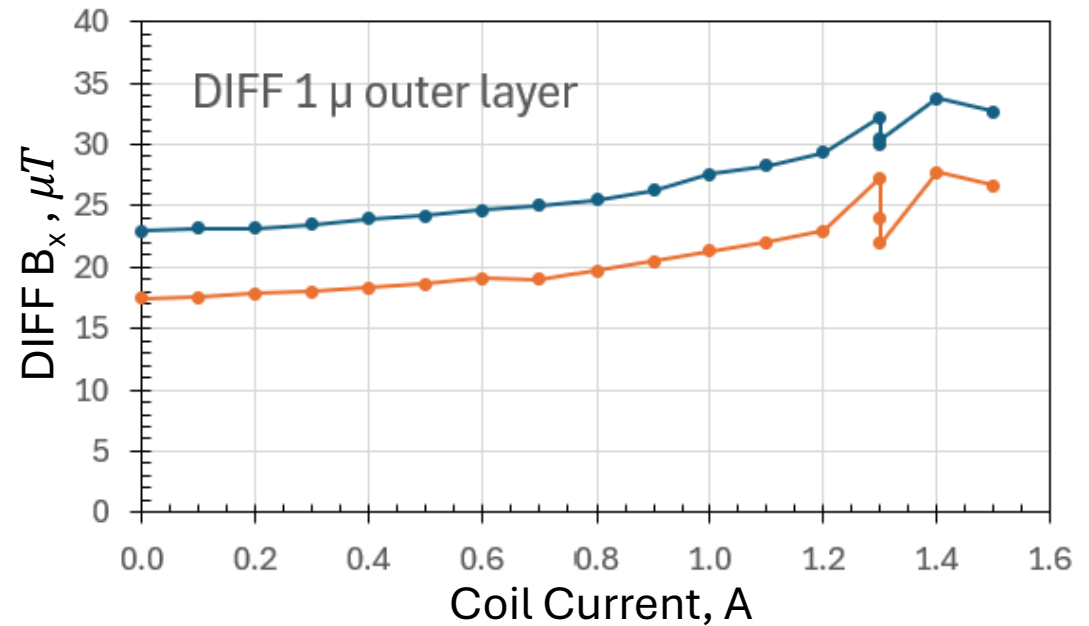
$$\sim 1.2965i^2 - 1.2367i + 3.5084 \mu T$$

# 1 $\mu$ - OUTER layers

Standard degaussing,  $B_x$  field with/without perturbation at z peak position 88"



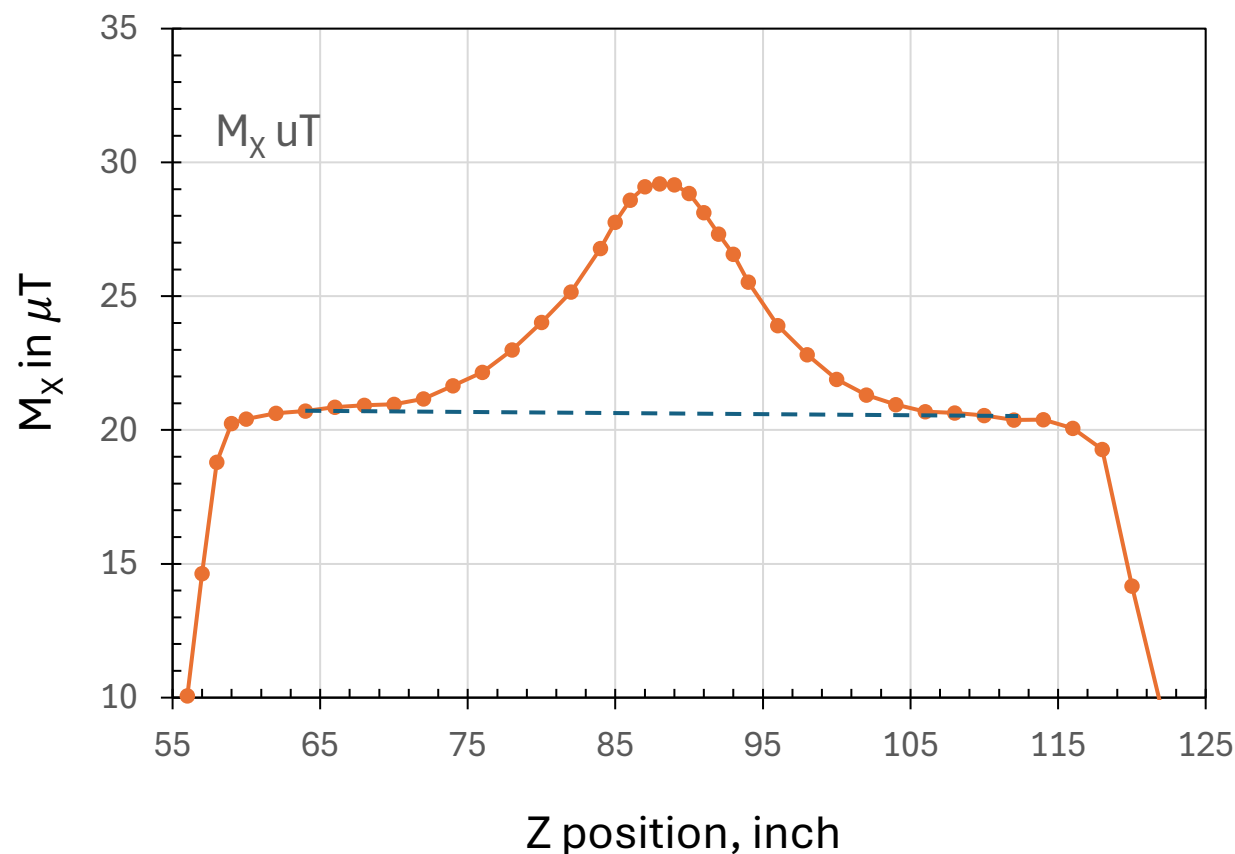
Single outer mu-layer  
seen PERT at 1.3 A  
 $\overline{\text{PERT}} = 27.6 \pm 3.7 \mu T$



DIFF dependence on current  
 $\sim 5.27i^2 - 1.26i + 20.46 \mu T$

## 2 $\mu$ - OUTER layers

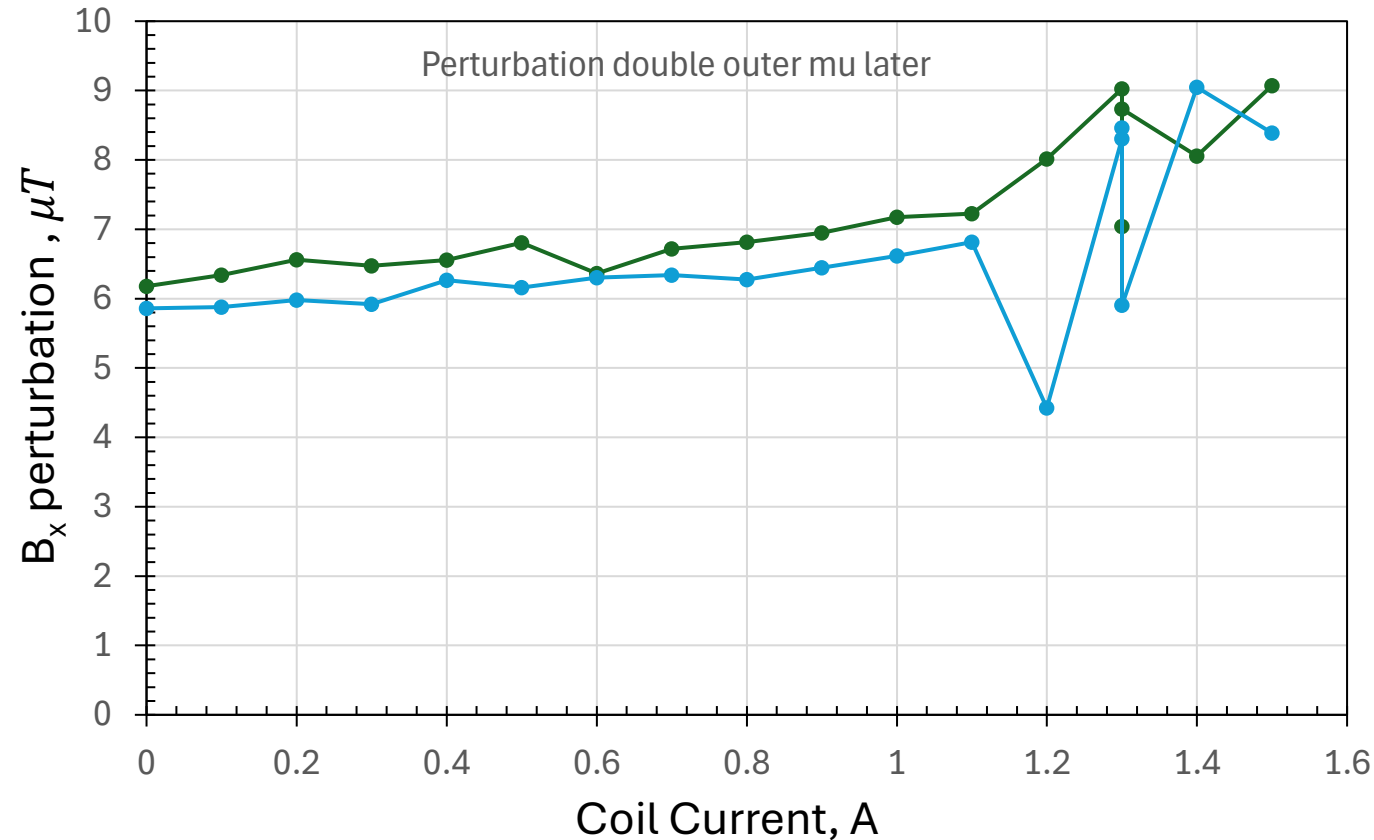
Special degaussing with 0A and NO PERT. Then, Z-scan with PERT, current 0 A



Peak position 88"  
Perturbation  $\sim 8.6 \mu T$

# 2 $\mu$ - OUTER layers

After Standard Degaussing, Z=88". DIFF=PERT-NOPERT, with Air



Double outer mu-layer  
Seen PERT at 1.3 A  
 $\overline{\text{PERT}} = 7.9 \pm 1.1 \mu T$

DIFF dependence on current  
 $\sim 1.713i^2 - 1.088i + 6.311 \mu T$

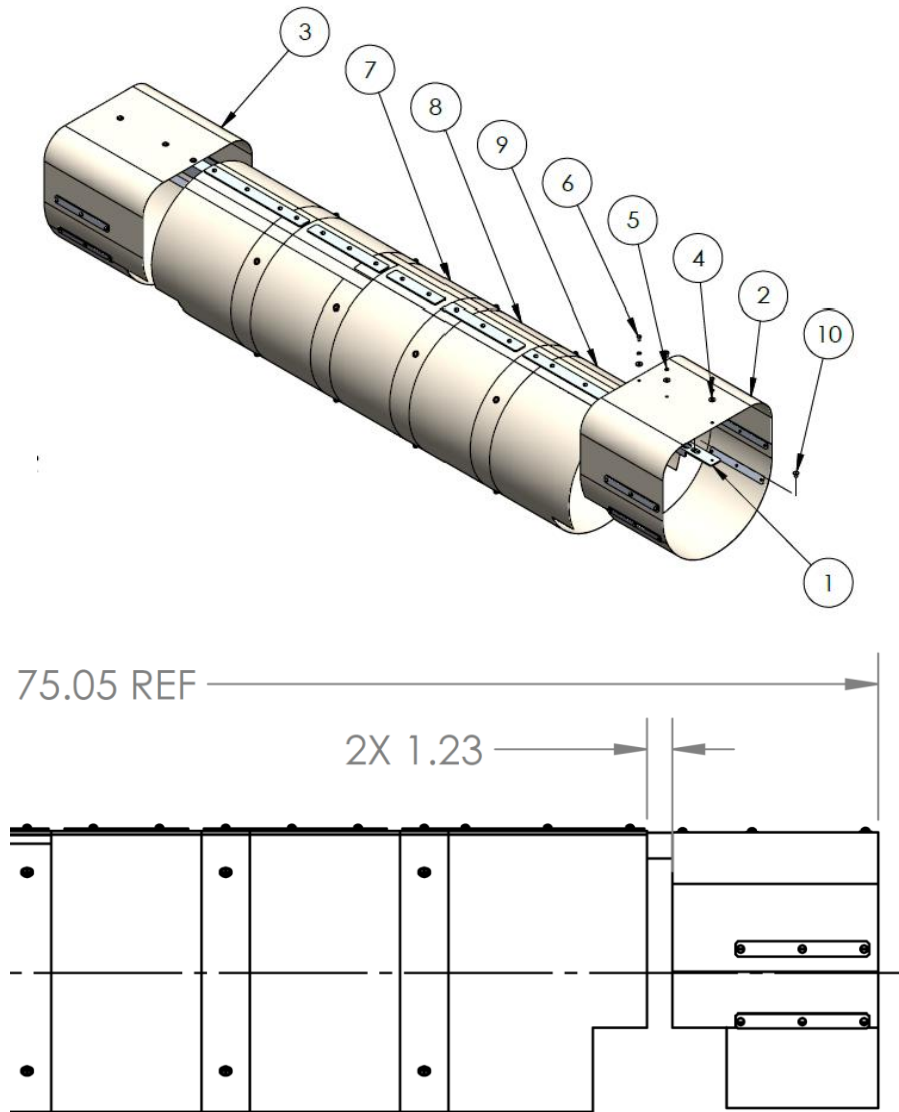
# Summary of shielding factors

All measurements are after standard degaussing with coil current 1.3 A (~ 25 Gauss)

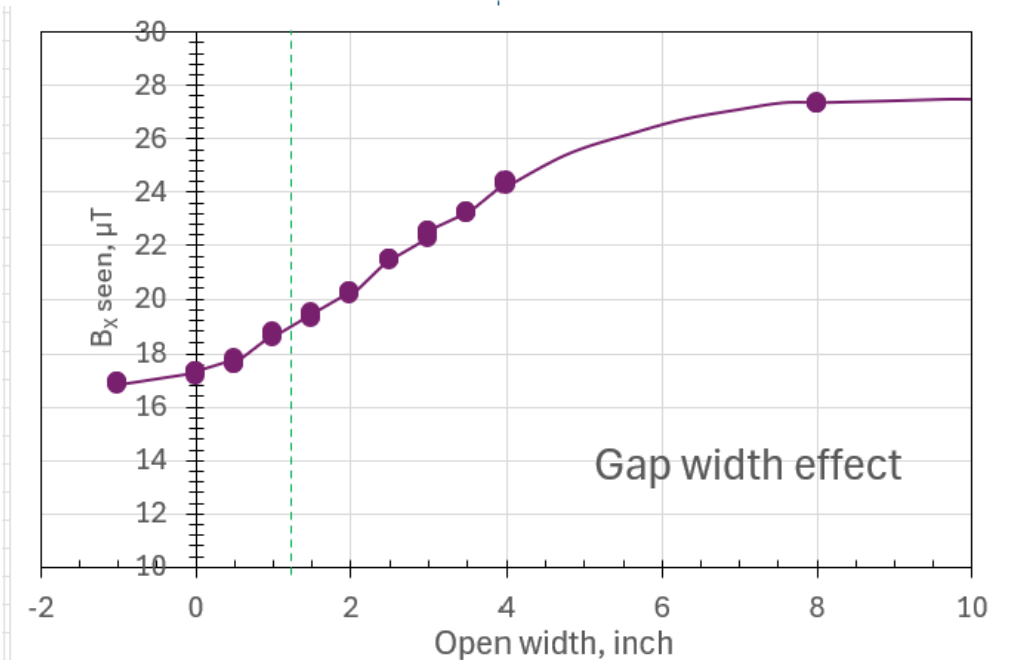
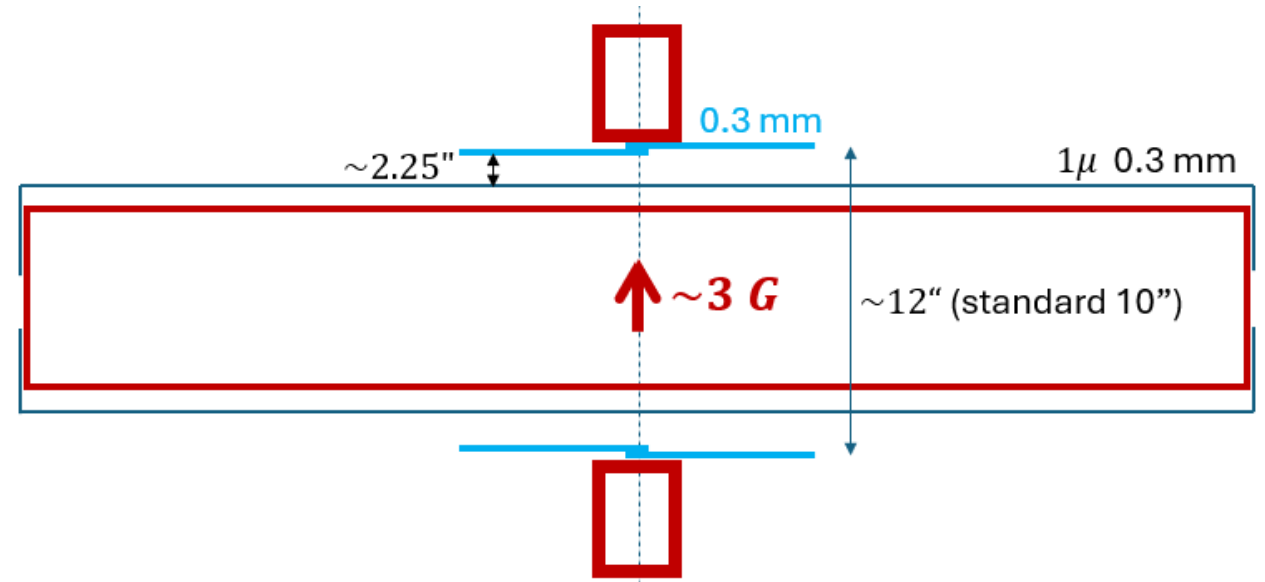
| Configuration       | Diameter, inch | Seen PERT $\mu T$ | SF               | $\mu_{eff}$ |
|---------------------|----------------|-------------------|------------------|-------------|
| 1 $\mu$ inner layer | 7.5            | 16.1 $\pm$ 1.8    | 32.6 $\pm$ 3.8   | 13,317      |
| 2 $\mu$ inner layer | 7.6            | 4.51 $\pm$ 0.61   | 116.2 $\pm$ 16.1 | 24,050      |
| 1 $\mu$ outer layer | 9.0            | 27.6 $\pm$ 3.7    | 19.0 $\pm$ 2.6   | 11,220      |
| 2 $\mu$ outer layer | 9.1            | 7.9 $\pm$ 1.1     | 66.3 $\pm$ 9.4   | 19,794      |

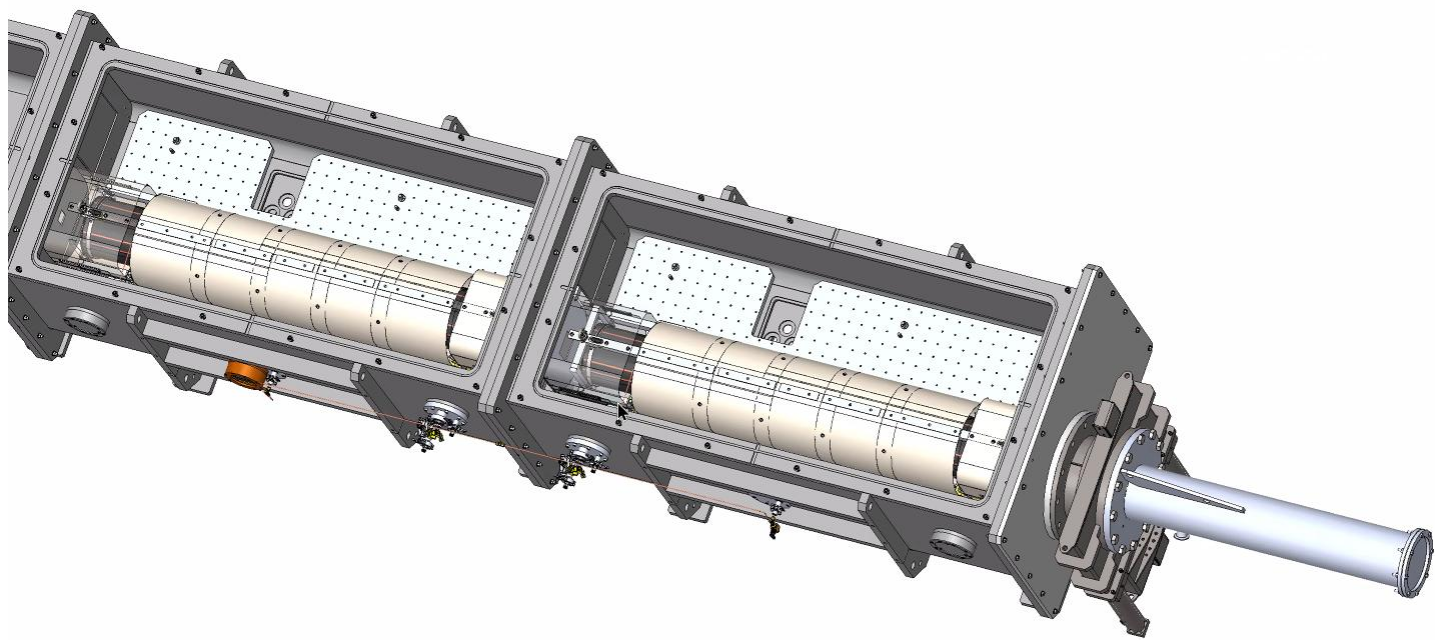
- We do not have degaussing of the outer layer – maybe we should
- Double layer seems have higher  $\mu$  probably due to effective space between two layers
- Layer thickness calculated effectively accounting for sheets overlay

# Gap in nTMM outer shield



# Gap simulated in prototype





# Changes From Previous Run

- CO<sub>2</sub> cans to be connected to each other (via 1/8" copper tubing) to ensure consistent CO<sub>2</sub> pressure
- CO<sub>2</sub> charge pressure to be increased 10% to mitigate any leakage during transportation or installation
- Preparation for the possibility of CO<sub>2</sub> gas recharge
- An additional borated poly beam stop is added between the CO<sub>2</sub> cans upstream of the cadmium beam stop
- The cadmium beam stop<sup>I</sup> previously located in the downstream end of chamber 7 is moved to the upstream end of chamber 8
- The outer magnetic shield has been augmented with a second layer of mu-metal for improved shielding + gap patch?
- Degaussing capability added to for outer shield